

Outflows from Accreting Compact Objects

Motivation

- Angular momentum loss
- "Veiling"
- AGN Unification
- AGN Unification
- Feedback
- State Transitions

Disk Winds vs Jets

	Jets	Disks Winds
Key Signature	Radio (sometimes resolved)	Blueshifted Absorption / P Cygni profiles
Speed	$v_j \sim c$ (relativistic)	$v_w \sim v_{esc}(R_{w,min})$ $\sim \max[v_{esc}(R_0), 0.1c]$
Geometry	highly collimated	weakly collimated -- bipolar [e.g. AWDs] -- equatorial [e.g. AGN]
Driving Mechanisms	(i) Blandford-Znajek / Penrose process (ii) MHD/Blandford-Payne (iii) Poynting Flux / Magnetic Tower	(i) MHD/Blandford-Payne (ii) Irradiation/Thermal Driving -- $v_{th} \gtrsim v_{esc}$ -- $T_{max} \simeq T_{Comp}$ -- $R_{Comp} = R[v_{th}(T_{Comp}) = v_{esc}]$ -- Effective @ $R \gtrsim 0.1R_{Comp}$ (iii) Radiation Pressure -- electron scattering ++ continuum driving ++ $L/L_{Edd} \sim 1$ -- line scattering ++ <i>line-driving</i> ++ $L/L_{Edd} \gtrsim 1/2000$
"Preferred" State	Hard State	Soft State
Mass-Loss Rate	?	$\dot{M}_j \sim 0.1 \dot{M}_{acc}$ (AWDs) $\dot{M}_w \sim 1 - 10 \dot{M}_{acc}$ (XRBS, AGN)
Kinetic Power	$\dot{M}_j v_j^2 \sim L_{acc} \sim GM\dot{M}_{acc}/R_0$ (bubbles as calorimeters)	$\dot{M}_w v_w^2 \sim 0.1L_{acc}$ (?)
Angular Momentum Lever Arm	$l_j \sim ???$	$l_w \sim ???$

Formation of P Cygni Profiles

- Monotonically Accelerating Spherical Outflows from Spherical Objects
 - (strictly) blue-shifted absorption
 - (strictly) symmetric emission
 - $v_{max} \simeq v_{\infty}$
 - $F_{max} @ v = 0$
 - pure scattering: $\int_{\lambda} f_{\lambda} d\lambda = 0$
 - photons encounter exactly **one** resonance
 - $\lambda_{photon} - \lambda_{line} = v/c$
 - defines *Sobolev region* where scattering is possible
 - all parts of the outflow are moving away from each other
- Rotating Disk Winds from Thin Disks
 - foreshortening
 - limb-darkening
 - wind geometry
 - wind kinematics
 - red-shifted absorption possible
 - for bipolar winds:
 - face-on: pure absorption
 - edge-on: pure emission
 - in general:
 - absorption when bright continuum is *behind* the wind
 - emission when it is not
 - pure scattering: $\int_{\Omega} \int_{\lambda} f_{\lambda} d\lambda d\Omega = 0$
 - **But**
 - collisional excitation → radiative de-excitation ("thermal emission")
 - collisional de-excitation ("thermalization")
 - recombination
 - photons may encounter **multiple** resonances

Thermally-Driven Winds

- For sufficiently strong (X-ray) irradiation, there will "always" be a layer in the disk atmosphere where Compton heating and cooling dominate
- This sets the maximum temperature in the atmosphere, $T_{max} \simeq T_{Comp} \simeq h \langle \nu \rangle / 4k$ Depends only on (local) SED
- The corresponding thermal speed is $v_{th}(T_{Comp}) \simeq \sqrt{kT_{Comp}/m_p}$
 - this is independent of radius (if $\tau \ll 1$)
- The escape velocity from the disk is $v_{esc} = \sqrt{2GM/R}$
- So beyond some critical radius -- the Compton radius -- the thermal speed exceeds v_{esc}
 - $R_{Comp} = R[v_{th}(T_{Comp}) = v_{esc}]$
- Thermal driving is pretty much inevitable at some level
- Solar wind is basically thermally driven
 - although it also has a significant magnetic lever arm, $R_A \simeq 20R_{\odot}$
 - **different driving mechanisms can act in concert**
- Properties of thermally driven winds depend critically on
 - *where* (radially and vertically) $v_{th} \simeq v_{esc}$
 - thermal stability
 - \rightarrow heating and cooling
 - local SED
 - optical depth / radiative transfer

Line-Driven Winds

- Eddington limit:
 - luminosity beyond which radiation pressure *on electrons* exceeds gravitational attraction
- But in many situations, $\sigma_e \ll \sigma_{b-b}$
 - line scattering can dominate the momentum transfer!
- Define "force multiplier" \mathcal{M} so that
 - $P_{rad,tot} = \mathcal{M}P_e$
 - **A magic number:** $\mathcal{M} \lesssim 2000$ (Gayley 1995)
- Line-driving is possible for $L \gtrsim L_{Edd}/2000$
- Driving can be dominated by:
 - a few strong resonance lines
 - many weaker lines
- Single-scattering limit:
 - Momentum carried by photons: L/c
 - Momentum carried by outflow: $\dot{M}_w v_\infty$
 - If each photon scatters once and transfer all of its *momentum* to the flow
 - $\dot{M}_{w,m} v_\infty \simeq L/c$
- Some line-driven flows (e.g. WR stars) violate this!
- Energy limit:
 - Energy carried by photons: L
 - Energy carried by outflow $\frac{1}{2}\dot{M}_w v_\infty^2$
 - If each photon transfers all of its *energy* to the flow
 - $\dot{M}_{w,e} v_\infty^2 \simeq 2L$
 - Requires that photons scatter many times
- Maximum achievable mass-loss rates:
 - $\frac{\dot{M}_{w,e}}{\dot{M}_{w,m}} \simeq 2 \frac{c}{v_\infty}$
- Properties depend critically on:
 - ionization structure
 - radiative transfer

Evidence for MHD/Centrifugal Driving in a BH XRB?

Miller+06

- Chandra X-ray spectra of GRO J1644:
 - eclipsing system:
 - $P_{orb} \simeq 2.6 d$
 - $i \simeq 76^\circ \pm 9^\circ$
 - $M_{BH} \simeq 7 M_\odot$
 - $M_2 \simeq 2.3 M_\odot$
 - $L_X \simeq 0.04 L_{Edd}$
 - BB Disk:
 - $kT \simeq 1.3 keV$
 - $L_{disk} \simeq 0.65 L_X$
 - PL Corona:
 - $\Gamma \simeq 3.5$
 - $L_{cor} \simeq 0.35 L_X$
 - $\simeq 75$ resonance lines showing blue-shifted absorption
 - $v_{shift} \simeq 300 - 1600 km/s$
 - high ionization lines
 - Fe XXIV (IP: 2.0 keV)
 - Si XiV (IP: 2.7 keV)
 - Mg XII (IP: 2.0 keV)
 - Photo-ionization modelling
 - constant density slab
 - $n \simeq 6e15 cm^3$
 - $\Delta R \simeq 2.5 \times 10^8 cm$
 - $R_w \simeq 4.8 \times 10^8 cm \simeq 200 R_{Schw}$
 - Inferred mass-loss rate
 - $\dot{M}_w \simeq 3.5 \times 10^{17} g/s \simeq \dot{M}_{acc}$
 - assumes $\Omega/4\pi \simeq 0.2$
 - $L_w \simeq \frac{1}{2} \dot{M}_w v_w^2 \simeq 4 \times 10^{32} erg/s \ll L_{acc}$
 - Ionization parameter:
 - $\log \xi = \log L_X / n R^2 \simeq 4.5 \pm 0.2$
 - Limits on location
 - $R_{w,min} \simeq 10^{7.5} cm$
 - otherwise n high enough to populate metastable FeXXIII levels
 - corresponding lines not observed
 - $R_{w,max} \simeq 10^{9.5} cm$
 - otherwise $\Delta R \gtrsim R_w$
 - dilution of radiation field through slab
 - cannot get high enough N at high enough ξ
 - Exclude thermal driving
 - $R_{Comp} \simeq 7 \times 10^{12} cm \gg R_w$
 - (actually $R_{Comp} \simeq 5 \times 10^{11} cm$, see below)
 - Exclude line driving
 - wind is too highly ionized
 - photo-ionization model gives $\mathcal{M} \simeq 2$

Are the Miller+06 arguments robust?

- Netzer06
 - can also fit the spectra with
 - $\log \xi \simeq 3$
 - $n \simeq 10^{13} \text{ cm}^{-3}$
 - $R \simeq 10^5 R_g \gg R_{Comp}$
 - thermal driving might work
- Miller+08, Kallman+09
 - Netzer models are a poorer fit than Miller models
 - Netzer models assume/require FeXXIV lines to be optically thick and partially covered
 - data suggests otherwise
 - **BUT** (Kallman+09)
 - $v_w \ll v_K(R_w)$ [for R_w in Miller's model]
 - $v_w \simeq v_K(R_{Comp})$
 - wide range in v_w
- Luketic+10
 - hydro simulations
 - difficult to produce thermally driven wind matching observations
- **General Point:**
 - **How much should one trust *any* one-zone photo-ionization model for a disk wind?**
 - Is there such a thing as "the" photo-ionization parameter?
 - Is there such a thing as "the" density?
 - ...
 - Luketic+10 models get closer to a realistic test
 - Over to Nick!